

## Index of refraction

$$n_r = \sqrt{1 + 4\pi n_{at} \alpha} \approx 1 + 2\pi n_{at} \alpha$$

$$\alpha = -\frac{e^2}{\hbar} |\langle \lambda | z | g \rangle|^2 \frac{1}{\delta}$$

neglecting counterrotating term,  $\delta = \omega - \omega_{kg}$   
single strong resonance

Def.  $\Gamma \equiv \frac{4}{3} \frac{\omega^3 e^2}{\hbar c^3} |\langle \lambda | z | g \rangle|^2$  natural linewidth

$$\Rightarrow n_r = 1 - 2\pi n_{at} \frac{3\Gamma}{4} \left( \frac{c^3}{\omega^3} \right) \frac{1}{\delta}$$

$$= \frac{\lambda^3}{(2\pi)^3}$$

$$= 1 - n_{at} \sigma_0 \frac{\lambda}{4\pi} \frac{\Gamma}{2\delta}$$

with def.  $\sigma_0 \equiv 6\pi \left( \frac{\lambda}{2\pi} \right)^2$  resonant cross section

Damping: now replace  $\delta \rightarrow \delta + i\gamma/2$

Def.  $\delta' = \delta / (\gamma/2)$

$$-\frac{\Gamma}{2\delta} \rightarrow \frac{\Gamma}{\gamma} \left( \frac{-1}{\delta' + i} \right) = \frac{\Gamma}{\gamma} \left[ \frac{i}{1 + \delta'^2} - \frac{\delta'}{1 + \delta'^2} \right]$$

$$e^{ikz} = e^{in_r \frac{\omega}{c} z} = e^{-\left[ \frac{\tilde{D}_0}{2} \left[ \frac{1}{1 + \delta'^2} + \frac{i\delta'}{1 + \delta'^2} \right] \right] z}$$

Dispersion      Absorption

optical density on resonance ( $\delta' = 0$ ):  $\tilde{D}_0 = n_{at} \sigma_0 \frac{\Gamma}{\gamma} z$

Note: When linewidth  $\gamma$  determined by spar. em  $k \rightarrow g$ , then  $\gamma = \Gamma$

Phase shift  $\phi$ :  $e^{ikz} = e^{-\frac{\tilde{D}_0}{2} \frac{1}{1 + \delta'^2}} e^{i\phi}$

$\phi = -\frac{\tilde{D}_0}{2} \frac{\delta'}{1 + \delta'^2}$ ; at  $\delta' = \pm 1$ : max phase shift  $\phi = \mp \frac{\tilde{D}_0}{4}$